

Topological string theory from Landau-Ginzburg models

based on: arXiv:0904.0862 [hep-th],
arXiv:1104.5438 & 1111.1749 [hep-th] with **Michael Kay**

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Outline

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 - A_∞ -algebras and relation to amplitudes
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 - bulk-deformed amplitudes \iff *curved* Calabi-Yau A_∞ -algebra
 - solution to deformation problem:
 - “weak” deformation quantisation
 - homological perturbation
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- open topological string theory \iff Calabi-Yau A_∞ -algebra
 - A_∞ -algebras and relation to amplitudes
 - B-twisted Landau-Ginzburg models
 - bulk-deformed amplitudes \iff *curved* Calabi-Yau A_∞ -algebra
 - solution to deformation problem:
 - “weak” deformation quantisation
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 - focus on general, conceptual results
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Open topological string theory

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Topological *field* theory **correlators**

$$\langle \psi_{i_1} \cdots \psi_{i_n} \rangle_{\text{disk}}, \quad \omega_{ij} = \langle \psi_i \psi_j \rangle_{\text{disk}}$$



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$$\left\langle \psi_{i_1} \psi_{i_2} \psi_{i_3} \int \psi_{i_4}^{(1)} \cdots \int \psi_{i_n}^{(1)} \right\rangle_{\text{disk}}, \quad \psi_i^{(1)} = [G_{-1}, \psi_i] d\tau$$

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$$W_{i_1 \dots i_n} = \left\langle \psi_{i_1} \psi_{i_2} \psi_{i_3} \int \psi_{i_4}^{(1)} \dots \int \psi_{i_n}^{(1)} \right\rangle$$

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Get **effective superpotential** from amplitudes:

$$\mathcal{W}(u) = \sum_{n \geq 3} \frac{1}{n} W_{i_1 \dots i_n} u_{i_1} \dots u_{i_n}$$

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$$\sum_{r,s} \pm \omega^{ii'} W_{i' i_1 \dots i_r j i_{r+s+1} \dots i_n} \omega^{jj'} W_{j' i_{r+1} \dots i_{r+s}} = 0$$

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A_∞ -algebras

An A_∞ -algebra is a graded vector space A together with a degree-one codifferential

$$\partial : T_A \longrightarrow T_A, \quad T_A = \bigoplus_{n \geq 1} A[1]^{\otimes n}, \quad \partial^2 = 0$$

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subject to the relations (from $\partial^2 = 0$)

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An **A_∞ -algebra** is a graded vector space A together with linear maps $m_n : A[1]^{\otimes n} \longrightarrow A[1]$ of degree $+1$ for all $n \geq 1$ such that

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$$n = 1 : m_1 \circ m_1 = 0$$

$$n = 2 : m_1 \circ m_2 + m_2 \circ (m_1 \otimes \mathbb{1}) + m_2 \circ (\mathbb{1} \otimes m_1) = 0$$

$$n = 3 : m_2 \circ (m_2 \otimes \mathbb{1}) + m_2 \circ (\mathbb{1} \otimes m_2) \\ + m_1 \circ m_3 + m_3 \circ (m_1 \otimes \mathbb{1}^{\otimes 2} + \mathbb{1} \otimes m_1 \otimes \mathbb{1} + \mathbb{1}^{\otimes 2} \otimes m_1) = 0$$

$$n = 4 : \dots$$

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An A_∞ -Algebra is **Calabi-Yau** if it is minimal and cyclic with respect to a non-degenerate pairing.

Relation to open topological string theory

Underlying TFT data (*Frobenius algebra*):

H : space of states = BRST cohomology with basis $\{\psi_i\}$

$\langle \psi_{i_0} \dots \psi_{i_n} \rangle$: correlators computed from OPE and topological metric

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How to compute the products m_n ?

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Minimal model theorem. For any A_∞ -algebra (A, ∂) , its cohomology $H = H_{m_1}(A)$ has a minimal A_∞ -structure $\tilde{\partial}$

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Sketch of proof: Compute Feynman diagrams.

$$\tilde{m}_n(\psi_{i_1} \otimes \dots \otimes \psi_{i_n}) = \sum_{\text{trees with } n \text{ leaves}} (-1)^{\#G} \text{Diagram}$$

General picture

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topological string field theory

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topological string field theory = *off-shell* space A

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minimal
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General picture

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topological string theory = *on-shell* space H with higher A_∞ -products $\tilde{\partial}$

(Same for bulk sector if we replace “DG” by “DG Lie” and “ A_∞ ” by “ L_∞ ”.)

B-twisted affine Landau-Ginzburg models

$$Z = \int \mathcal{D}\Phi \, e^{-\int K(\Phi, \bar{\Phi}) - (\frac{i}{2} \int W(\Phi) + \text{c. c.})} \, \text{str} \left(\mathcal{P} \, e^{-\oint (\frac{1}{2} \rho^i \nabla_i D + \frac{i}{2} \eta^{\bar{i}} \nabla_{\bar{i}} D^\dagger + \frac{i}{2} \{D, D^\dagger\})} \right)$$

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topological metric:

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Defect sector. $\dots \Rightarrow$ extended 2d TFT

Open top. string theory for Landau-Ginzburg models

Apply minimal model theorem to off-shell algebra

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Result. General algorithm to construct all open tree-level amplitudes, and (another) first-principle derivation of Kapustin-Li pairing.

Bulk deformations

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Fact. Curvature screws everything up:

- m_1 no longer a differential
- cannot apply minimal model theorem
- need a new approach

Deformations and Maurer-Cartan equations

Given an A_∞ -algebra (A, ∂) , a *deformation* is $\delta \in \text{End}^1(T_A)$ such that $(A, \partial + \delta)$ is a curved A_∞ -algebra.

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Fact. Let L be an L_∞ -quasi-isomorphism between DG Lie algebras. Then

$$\delta \longmapsto \sum_{n \geq 1} \frac{1}{n!} L_n(\delta^{\wedge n})$$

is an *isomorphism* between the spaces of Maurer-Cartan solutions modulo gauge transformations.

Back to Landau-Ginzburg models

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Off-shell bulk sector is also a DG Lie algebra:

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Transport them to deformations of $(H, \tilde{\partial})$ via an L_{∞} -map

$$\left(T_{\text{poly}}, [-W, \cdot]_{\text{SN}}, [\cdot, \cdot]_{\text{SN}} \right) \longrightarrow \left(\text{Coder}(T_H), [\tilde{\partial}, \cdot], [\cdot, \cdot] \right)$$

Back to Landau-Ginzburg models

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First step: deformations of the off-shell open string algebra

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First step: deformations of the off-shell open string algebra

Theorem. There is a sequence of explicit L_∞ -quasi-isomorphisms

$$\begin{aligned} \left(T_{\text{poly}}, [-W, \cdot]_{\text{SN}}, [\cdot, \cdot]_{\text{SN}} \right) &\xrightarrow[\text{deform. quant.}]{\text{"weak"}} \left(\text{Coder}(T_{\mathbb{C}[\mathbf{x}]}) , [\hat{\partial}_0 + \hat{\partial}_2, \cdot], [\cdot, \cdot] \right) \\ &\xrightarrow[\text{equivalence}]{\text{Morita}} \left(\text{Coder}(T_A), [\partial_0 + \partial_2, \cdot], [\cdot, \cdot] \right) \\ &\xrightarrow[\text{cancellation}]{\text{tadpole}} \left(\text{Coder}(T_A), [\partial_1 + \partial_2, \cdot], [\cdot, \cdot] \right) \end{aligned}$$

Digression: deformation quantisation à la Kontsevich

Consider classical theory with phase space $M = \mathbb{R}^d$ and associative, commutative algebra of observables $(C^\infty(M, \mathbb{R}), \cdot) \equiv (C^\infty(M, \mathbb{R}), \hat{\partial}_2)$.

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Deformation quantisation constructs algebra of quantum observables $(C^\infty(M, \mathbb{R})[[\hbar]], \star)$ by deforming the product to

$$f \star g = f \cdot g + B_1(f, g)\hbar + B_2(f, g)\hbar^2 + \dots$$

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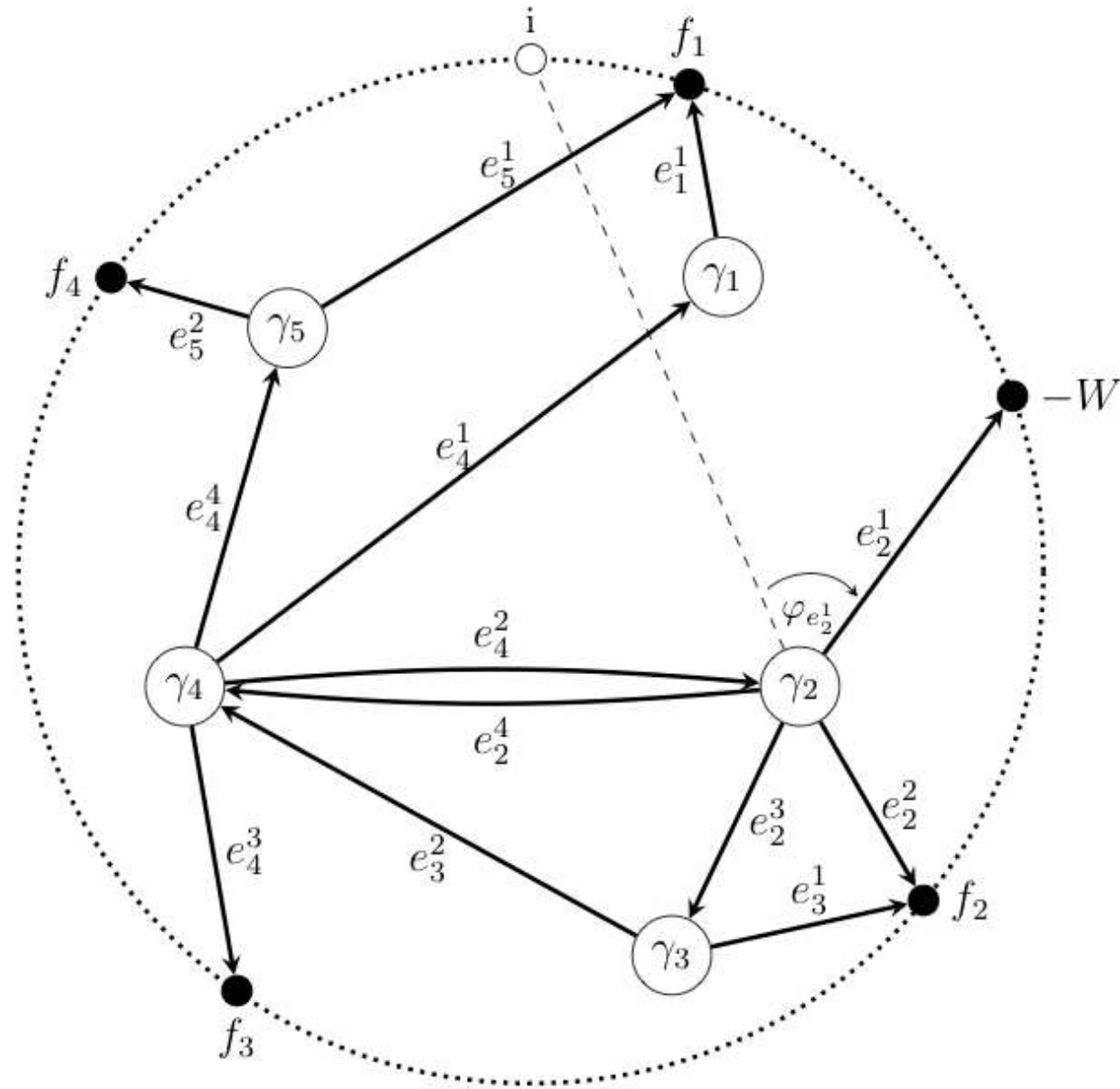
Kontsevich constructs an explicit L_∞ -quasi-isomorphism

$$K : \left(\Gamma(M, \bigwedge TM), 0, [\cdot, \cdot]_{\text{SN}} \right) \longrightarrow \left(\text{Coder}(T_{C^\infty(M, \mathbb{R})}), [\widehat{\partial}_2, \cdot], [\cdot, \cdot] \right)$$

Digression: deformation quantisation à la Kontsevich

$$\begin{aligned}
 & (K_n(\gamma_1 \wedge \dots \wedge \gamma_n))_m (f_1 \otimes \dots \otimes f_m) \\
 &= \sum_{\Gamma \in \mathcal{G}(n,m)} \frac{1}{(2\pi)^{2n+m-2}} \int_{\iota(\overline{C}^{n,m})} \bigwedge_{k=1}^n (d\varphi_{e_k^1} \wedge \dots \wedge d\varphi_{e_k^{\tilde{\gamma}_k}}) \\
 & \quad \cdot \sum_I \left[\prod_{i=1}^n \left(\prod_{e \in \Gamma_{\bullet \rightarrow i}} \partial_{I(e)} \right) \gamma_i^{I(e_1^1) \dots I(e_i^{\tilde{\gamma}_i})} \right] \left[\prod_{\bar{j}=1}^{\bar{m}} \left(\prod_{e \in \Gamma_{\bullet \rightarrow \bar{j}}} \partial_{I(e)} \right) f_{\bar{j}} \right]
 \end{aligned}$$

Digression: deformation quantisation à la Kontsevich



Weak deformation quantisation

Theorem. Kontsevich's map

$$K : \left(T_{\text{poly}}, 0, [\cdot, \cdot]_{\text{SN}} \right) \longrightarrow \left(\text{Coder}(T_{\mathbb{C}[\mathbf{x}]}) , [\widehat{\partial}_2, \cdot], [\cdot, \cdot] \right)$$

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is also an L_∞ -quasi-isomorphism

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Thus we have constructed the first part of our bulk deformation map

$$\begin{aligned} \left(T_{\text{poly}}, [-W, \cdot]_{\text{SN}}, [\cdot, \cdot]_{\text{SN}} \right) &\longrightarrow \left(\text{Coder}(T_A), [\partial, \cdot], [\cdot, \cdot] \right) \\ &\longrightarrow \left(\text{Coder}(T_H), [\tilde{\partial}, \cdot], [\cdot, \cdot] \right) \end{aligned}$$

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Theorem. Let (A, ∂) be an A_∞ -algebra and $(H, \tilde{\partial})$ its minimal model.

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$$(T_H, \tilde{\partial}) \begin{array}{c} \xrightarrow{F} \\ \xleftarrow{\bar{F}} \end{array} (T_A, \partial) \begin{array}{c} \circlearrowleft \\ \circlearrowright \end{array} U$$

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in standard form, i. e.

$$\bar{F}F = \mathbb{1}_{T_H}, \quad \mathbb{1}_{T_A} - F\bar{F} = \partial U + U\partial, \quad U^2 = UF = \bar{F}U = 0$$

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This gives rise to an L_∞ -morphism

$$\begin{aligned} \left(\text{Coder}(T_A), [\partial, \cdot], [\cdot, \cdot] \right) &\longrightarrow \left(\text{Coder}(T_H), [\tilde{\partial}, \cdot], [\cdot, \cdot] \right) \\ \delta &\longmapsto \sum_{n \geq 1} \bar{F}(\delta U)^n \delta F \end{aligned}$$

Recursive formulas for \bar{F}, U

$$U_n^1 = -\frac{1}{2}G\partial_2^1 \left(\sum_{l=1}^{n-1} (U_l^1 \otimes (\mathbb{1} + F\bar{F})_{n-l}^1 + (\mathbb{1} + F\bar{F})_{n-l}^1 \otimes U_l^1) \right)$$

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Bulk-deformed amplitudes for Landau-Ginzburg models

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All off-shell deformations δ are bulk-induced, i. e. uniquely defined by

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Thus the curved A_∞ -products describing bulk-deformed open topological string amplitudes are explicitly encoded in

$$\tilde{\partial} + \sum_{n \geq 1} \bar{F}(\delta U)^n \delta F$$

Conclusion

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- open topological string theory \iff Calabi-Yau A_∞ -algebra
- gives explicit algorithm to compute open amplitudes
- **bulk-deformed amplitudes** computable via *weak deformation quantisation* and *homological perturbation*:

$$\begin{aligned} & \left\langle \psi_{i_0}, \tilde{m}_n^t(\psi_{i_1} \otimes \dots \otimes \psi_{i_n}) \right\rangle \\ &= \left\langle \psi_{i_0}, \psi_{i_1} \psi_{i_2} \int \psi_{i_3}^{(1)} \dots \int \psi_{i_n}^{(1)} e^{\sum_i t_i \int \phi_i^{(2)}} \right\rangle \end{aligned}$$